

# Extension of susceptibilities, screened exchange and spin-fluctuation integrals into ultrasoft pseudopotentials

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In the Eliashberg theory or density functional theory for superconductors (SCDFT) [1], where the electron-phonon, screened Coulomb, spin-fluctuation mediated interactions are included non-empirically, we need to compute the following product of two Kohn-Sham orbitals

$$\rho_{n\mathbf{k}n'\mathbf{k}'}(\mathbf{r}) = \varphi_{n\mathbf{k}}^*(\mathbf{r})\varphi_{n'\mathbf{k}'}(\mathbf{r}), \quad (1)$$

where  $n$  ( $n'$ ) and  $\mathbf{k}$  ( $\mathbf{k}'$ ) are the band index and Bloch wavenumber, respectively. To perform this product together with the ultrasoft pseudopotentials (USPP) or projector augmented waves (PAW) that are widely used because of the good accuracy and reasonable numerical costs [2], we need a correction term for the norm-conservation. Such a correction is originally proposed for the calculation of the susceptibility [3] as

$$\Delta\rho_{n\mathbf{k}n'\mathbf{k}'}(\mathbf{r}) = \sum_{\tau ii'} \langle \varphi_{n\mathbf{k}} | \beta_{\tau i'} \rangle \langle \beta_{\tau i} | \varphi_{n'\mathbf{k}'} \rangle Q_{\tau ii'}(\mathbf{r}), \quad (2)$$

where  $\beta_{\tau i}$  is the projector dual to the atomic pseudo orbital  $\psi_{\tau i}^{PS}$  at atom  $\tau$  and orbital  $i$ , and the augmentation charge  $Q_{\tau ii'}(\mathbf{r})$  is computed from the pseudo (PS) and all-electron (AE) atomic orbitals as follows:

$$Q_{\tau ii'}(\mathbf{r}) \equiv \psi_{\tau i}^{AE*}(\mathbf{r})\psi_{\tau i'}^{AE}(\mathbf{r}) - \psi_{\tau i}^{PS*}(\mathbf{r})\psi_{\tau i'}^{PS}(\mathbf{r}). \quad (3)$$

To utilise this correction to the calculation of the spin-fluctuation mediated interaction,

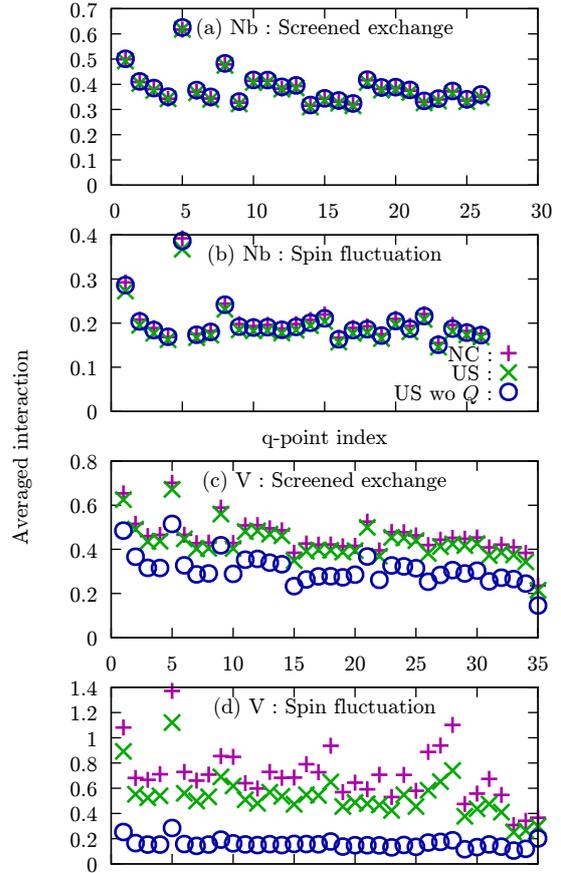


Figure 1: Spin-fluctuation mediated- (a, c) and screened exchange- (b, d) interactions averaged over the Fermi surface of Nb (a, b) and V (c, d). “+”, “×”, and “o” indicate results with norm-conserving (NC) ultrasoft (US) with and without augmentation charge  $Q(\mathbf{r})$ , respectively.

we implemented this formalism into our first-principles program package Superconducting-Toolkit [4] which is based on SCDFE.

Figure 1 shows the screened exchange and spin-fluctuation interactions averaged over the Fermi surface of Nb and V. In Nb, the effect of the augmentation charges  $Q$  is small because the USPP for this atom is almost norm-conserving. At the same time, V has a significant contribution from  $Q$  due to the nodeless 3d orbitals. In Fig. 1(d), we can see a deviation between the result by NC and US pseudopotential even if we include the correction. This deviation may be because the exchange-correlation kernel included in the spin-fluctuation is sensitive to the charge density in the vicinity of atoms.

We also performed the benchmark of the calculation of  $T_c$  for 15 materials, namely Al, V, Ta, In, Zn, Cd, Sn, ZrN, TaC, MgB<sub>2</sub>, H<sub>3</sub>S (at a pressure of 200 GPa), CaC<sub>6</sub>, YNi<sub>2</sub>B<sub>2</sub>C, and V<sub>3</sub>Si. Figure 2 shows the experimental, and calculated  $T_c$ ; we performed four kind of calculations by changing superconducting density functional, namely the conventional plasmon-assisted [5], Sanna’s Eliashberg-mimic (Sanna) [6], Sanna+Coulomb renormalization ( $Z_C$ ) [7], and Sanna +  $Z_C$  + Spin-fluctuation [9, 8] functional, respectively.

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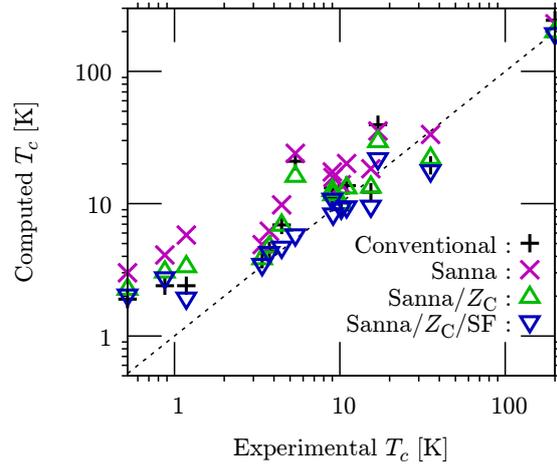


Figure 2: Computed- and experimental  $T_c$ . “+”, “x”, “ $\Delta$ ”, and “ $\nabla$ ” indicate result with conventional plasmon-assisted [5], Sanna’s Eliashberg-mimic (Sanna) [6], Sanna+Coulomb renormalization ( $Z_C$ ) [7], and Sanna +  $Z_C$  + Spin-fluctuation [9, 8] functional, respectively.

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